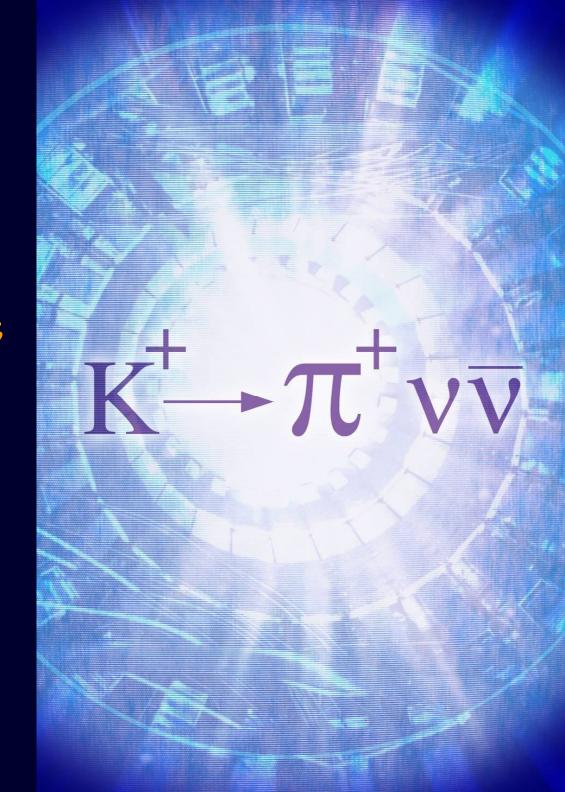
Final BNL E949 results on the measurement of the rare decay $K^+ \rightarrow \pi^+ \nu \bar{\nu}$

Joss Ives

University of the Fraser Valley



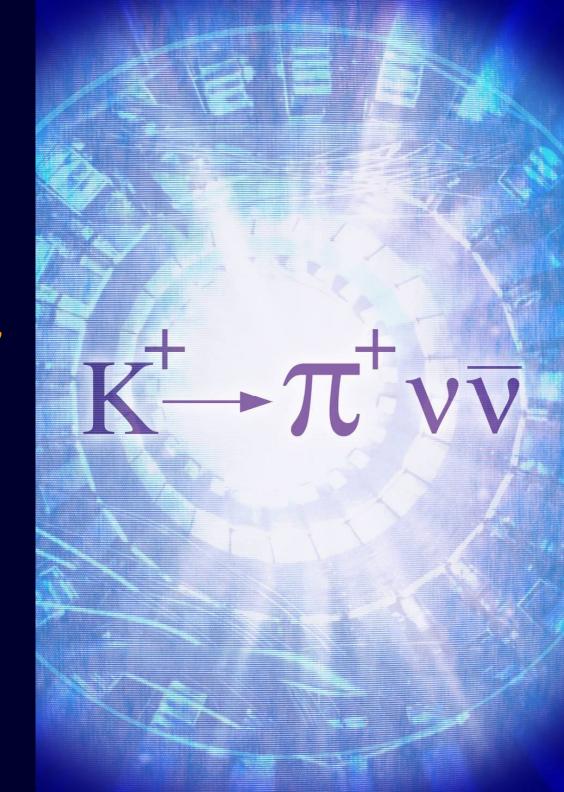
Final BNL E949 results on the measurement of the rare decay $K^+ \rightarrow \pi^+ \nu \bar{\nu}$

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UBC / TRIUMF







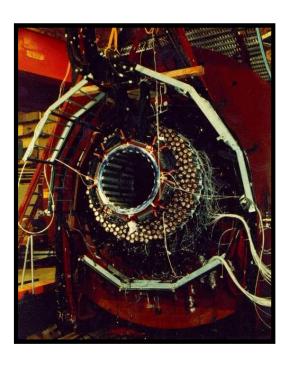
To find physics beyond the Standard Model, particle physicists try to break the Standard Model

Energy frontier



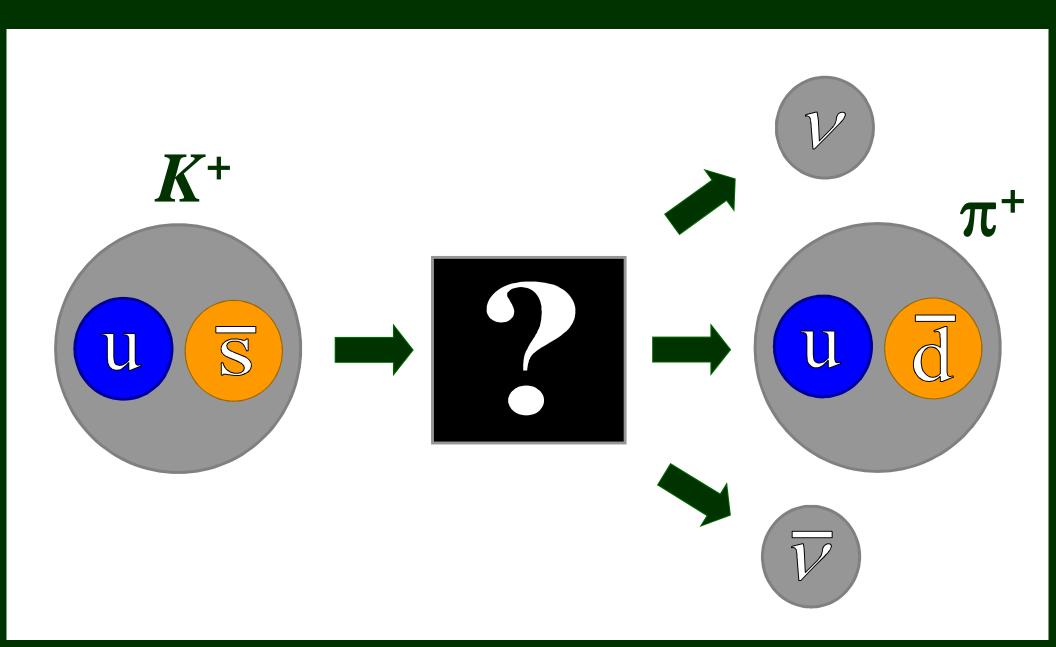
LHC

Precision Frontier

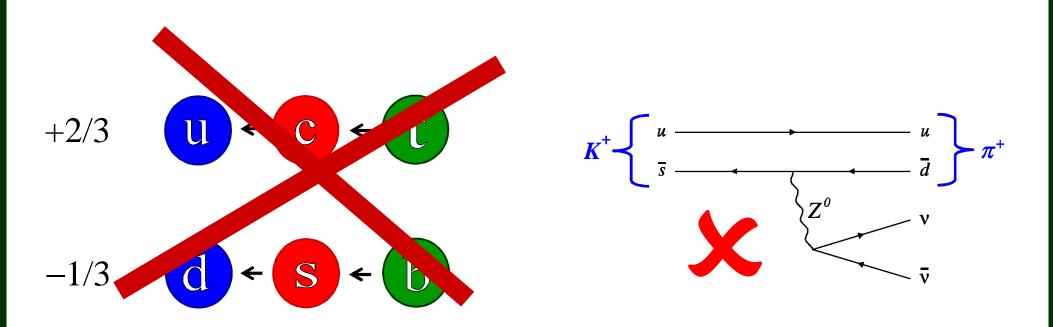


Rare decays such as $K^+ \rightarrow \pi^+ \nu \overline{\nu}$

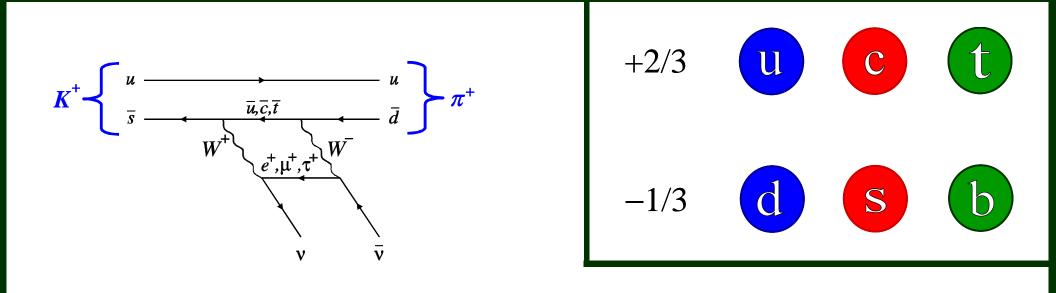
To understand why $K^+ \to \pi^+ \nu \bar{\nu}$ is rare, we need to look at the decay of s-bar to d-bar

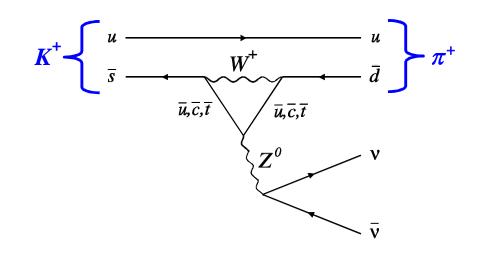


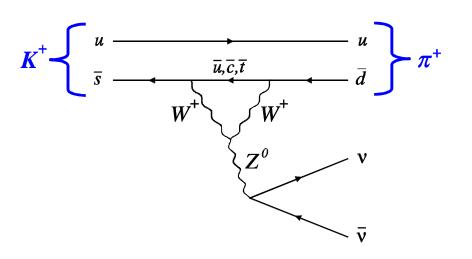
$K^+ \to \pi^+ v \overline{v}$ cannot proceed through a first-order weak process



But $K^+ \to \pi^+ \nu \overline{\nu}$ can proceed very slowly through second-order weak processes







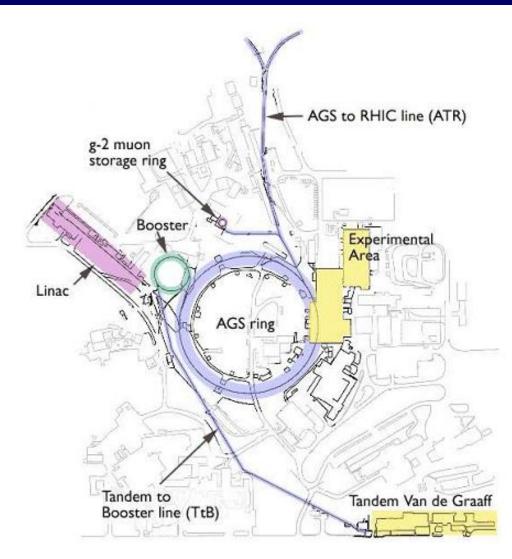
Observing the rare decay $K^+ \to \pi^+ \nu \bar{\nu}$ is an enormous experimental challenge

Standard Model branching ratio is $(0.85 \pm 0.7) \times 10^{-10}$

1 in 12.5 billion!

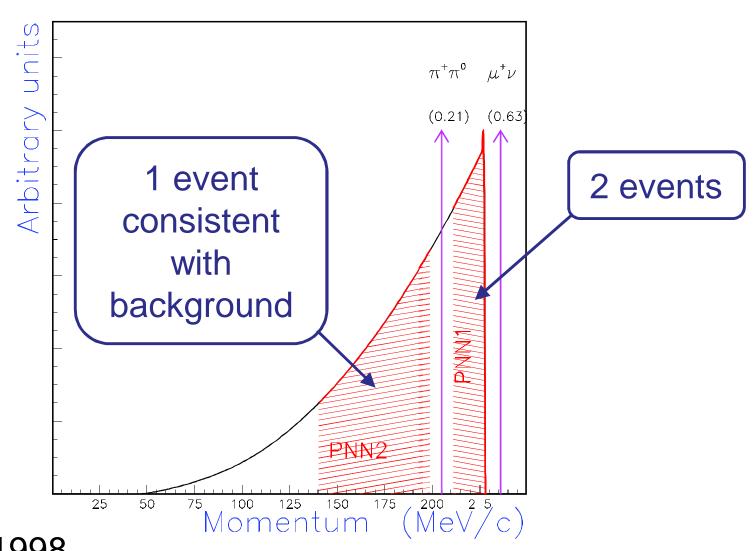
Experimental signature is π^+ + "nothing"

The first candidate events were observed by the E787 experiment at Brookhaven National Lab



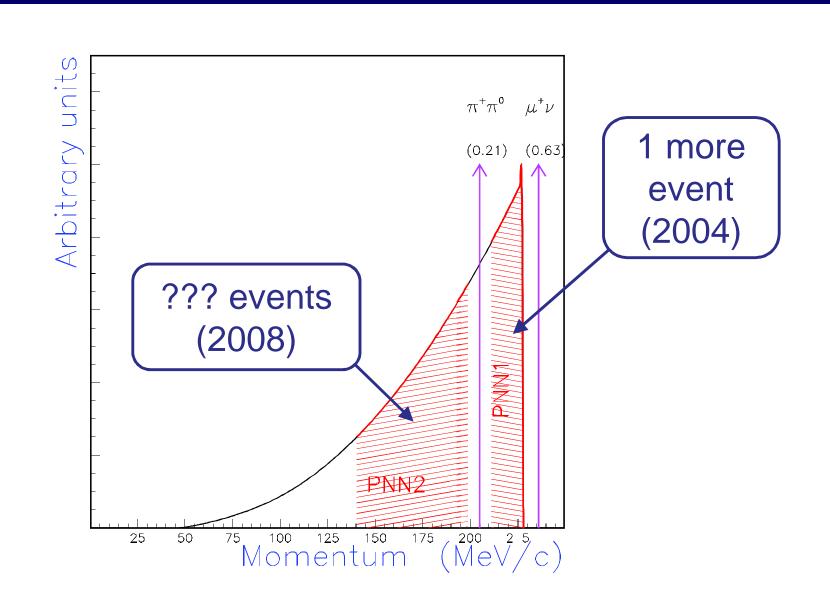
The Alternating Gradient Synchrotron (AGS)

The $K^+ \to \pi^+ \pi^0$ decay splits the momentum phase space in two, creating two regions of high signal-to-background

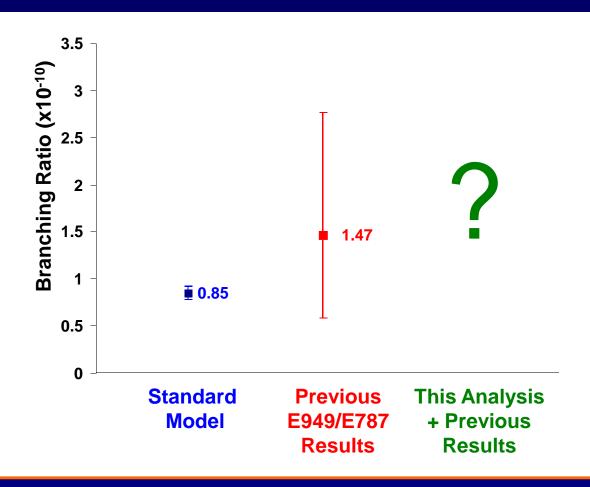


E787 = 1988 - 1998

The E949 experiment collected data in 2002 after major upgrades to the detector and beamline



BNL E949 & E787 have previously observed four candidate events showing that it can be done

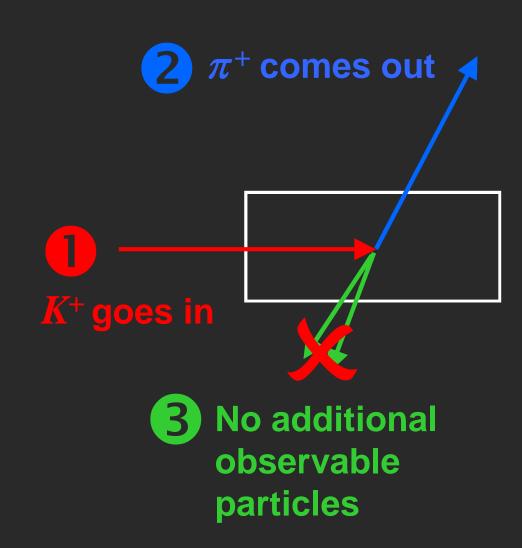


A precision measurement of $K^+ \to \pi^+ \nu \bar{\nu}$ probes for new physics

The rest of the talk focuses on the PNN2 region and how E949 overcame the difficulty of observing $K^+ \to \pi^+ \nu \bar{\nu}$

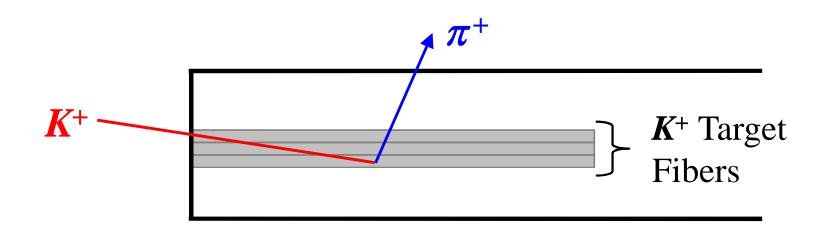
Experimental Techniques Analysis Techniques Final results

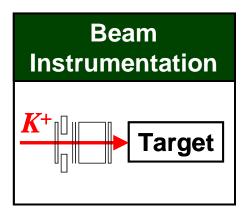
The experimental challenge is to observe π^+ + "nothing"



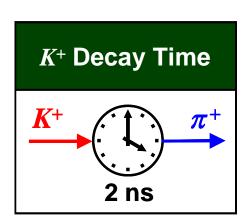


700 MeV/c kaon enters and decays-at-rest in the scintillator-fiber target





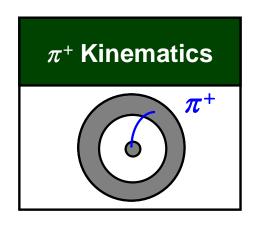
Beam instrumentation is used to identify a single K^+ from the beam

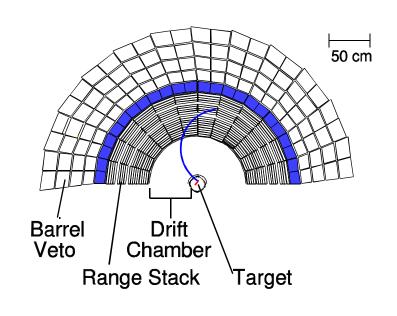


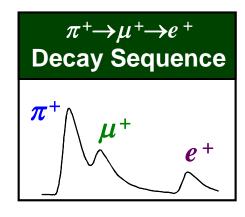
Kaon comes to rest in the target and we wait at least 2 ns for K^+ to decay



Pion comes to rest in the scintillator rangestack after traversing the drift chamber





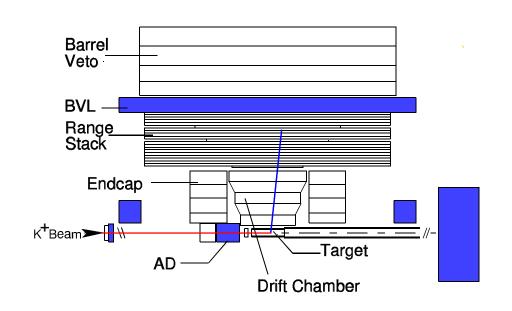


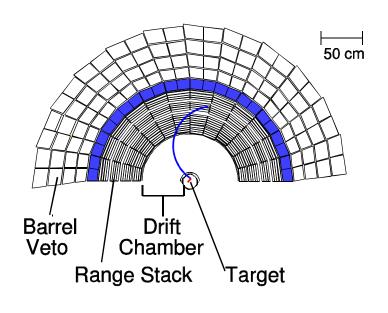
 π^+ momentum is measured in the drift chamber

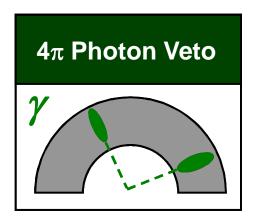
 π^+ range and energy are measured target and range-stack

3

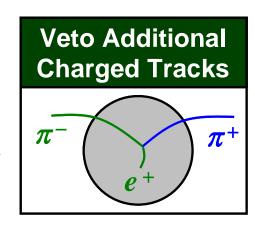
Veto additional charged particle and photon activity





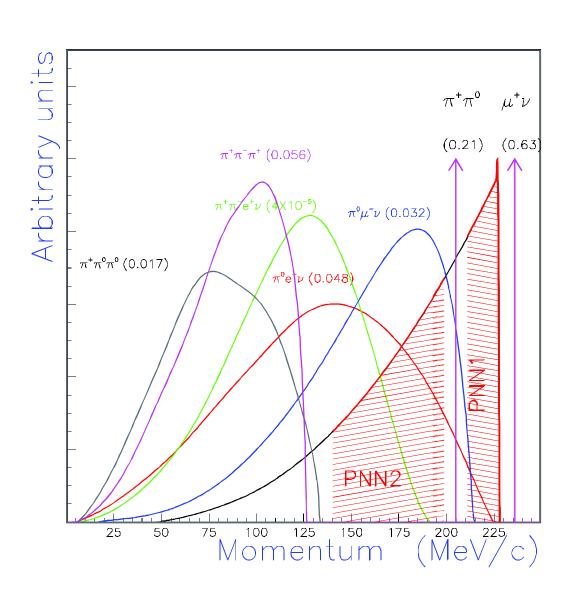


Pattern recognition is used in the target, drift chamber and range-stack



To reduce potential bias, we use a "blind" analysis technique

In our blind analysis, backgrounds that can mimic $K^+ \to \pi^+ \nu \overline{\nu}$ are identified *a priori*



Our blind analysis technique is based on eliminating everything that is not $K^+ \to \pi^+ \nu \overline{\nu}$ without directly examining our signal region

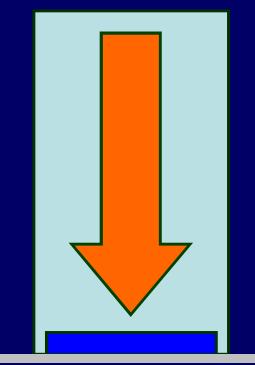
Take all events in the PNN2 high signal-to-background region

Events identified as being NOT $K^+ \rightarrow \pi^+ \nu \nu$

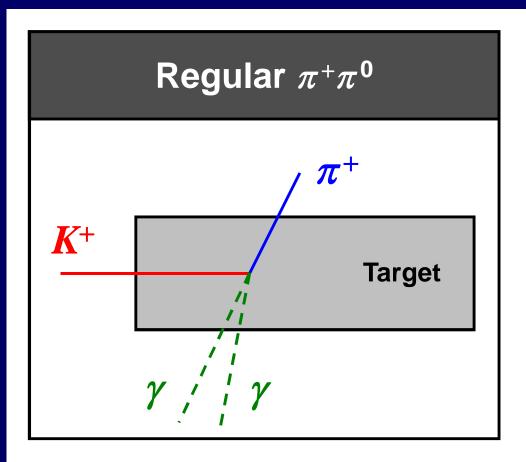
Candidate $K^+ \rightarrow \pi^+ \nu \nu$ events

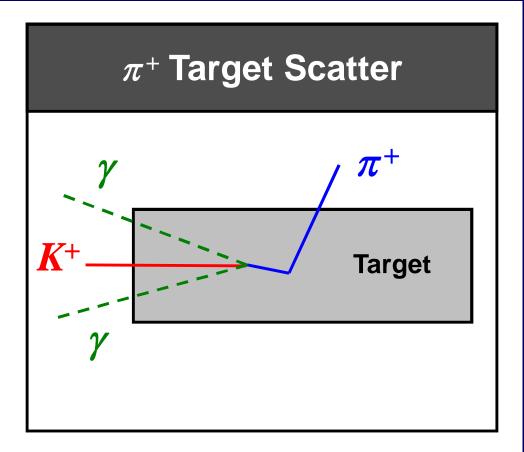
The key is VERY large suppression of background





The most troublesome background is $\pi^+\pi^0$ when the π^+ scatters in the target

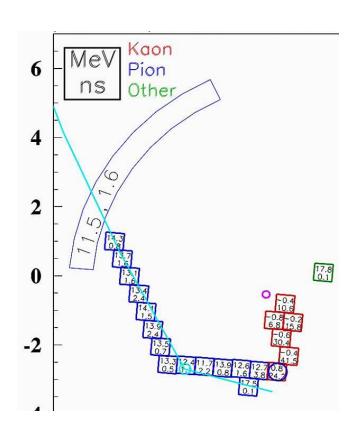




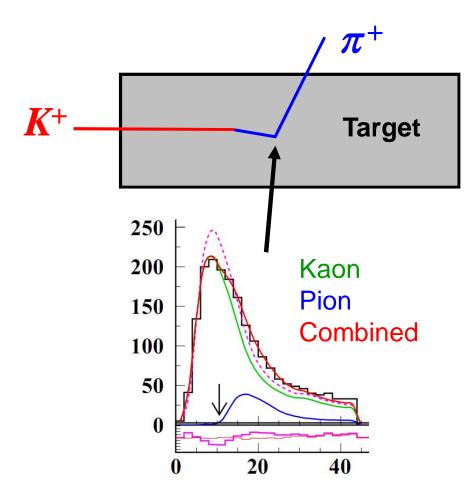
 π^+ and $\pi^{\,0}$ are back to back so photons are directed to efficient part of the photon veto

When the π^+ scatters in the target it loses energy and the photons lose directional correlation with the π^+

This background is suppressed by the photon veto and identification of π^+ scattering

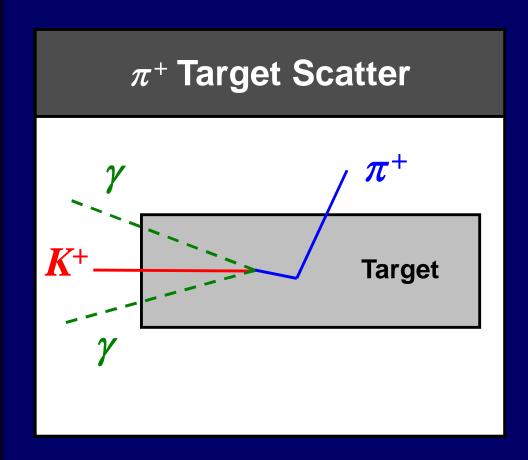


A scatter can be identified by finding a kink in the π^+ track



A scatter can also be identified by large energy deposits in a pion fiber or hidden energy in a kaon fiber

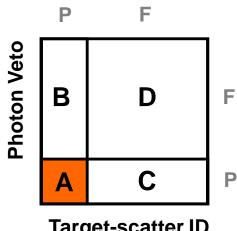
The π^+ target-scatter background will be used to illustrate the bifurcation background estimation technique



The photon veto and the target-scatter ID cuts are used to define the bifurcation phase space

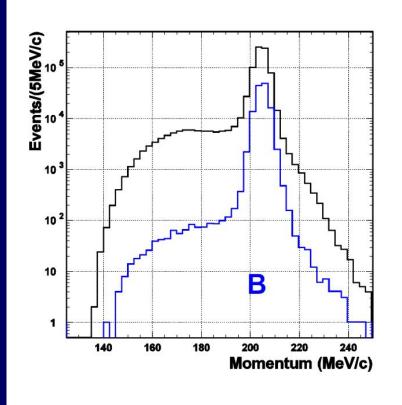
The signal region A cannot be examined directly. For uncorrelated cuts:

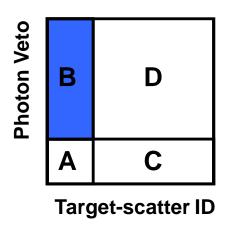
A/B = C/D



Target-scatter ID

Photon veto is inverted (black curve) and the target cuts are applied (blue curve)

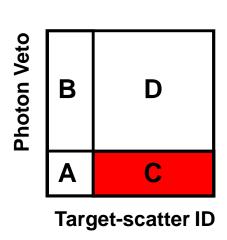


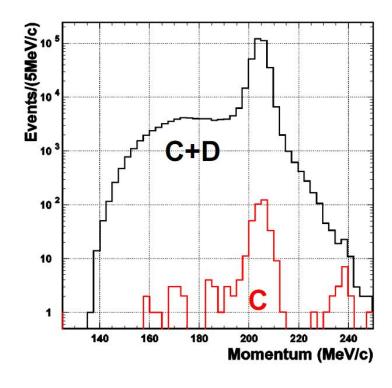


Measure B

Target-scatter cuts are inverted (black curve) and the photon veto is applied (red curve)

Measure C+D and C

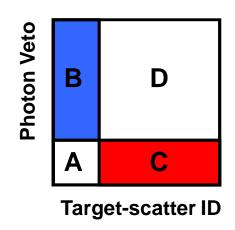




The number of π^+ target-scatter events is estimated using the number of events in each of the other regions

Background estimate:

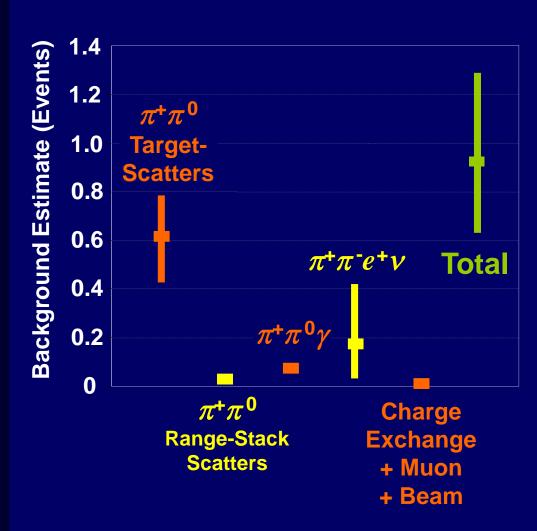
A = BC/D



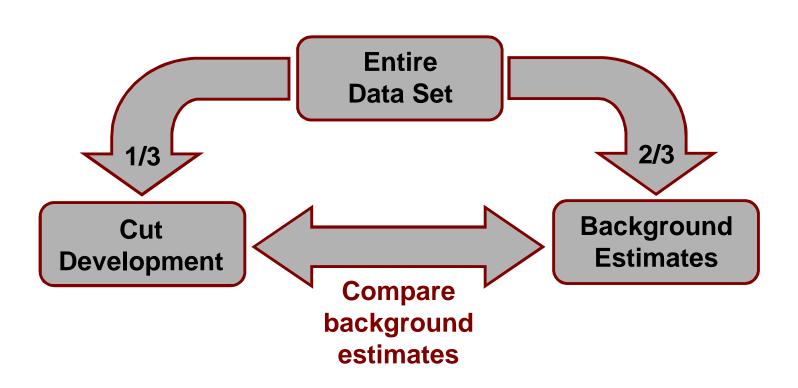
Result:

 $0.619 \pm 0.150 ^{+0.067}_{-0.100}$

The total background was suppressed to below a single event from 1.70 × 10¹² kaon decays with a total acceptance of 1.37 × 10⁻³



Background estimates are performed using data different from those used to develop cuts



Validation studies were performed...

To look for correlations between sets of bifurcation cuts

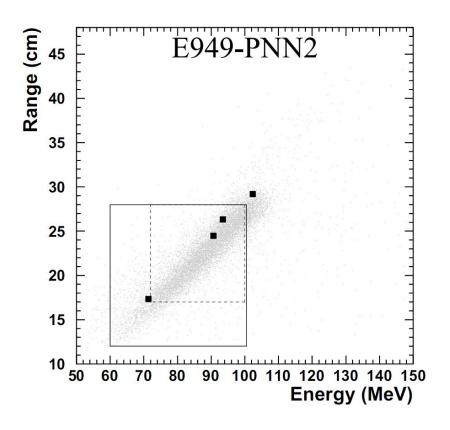
To search for additional background sources

The signal region is only examined once all background estimates have been finalized.

It's time to open the box...

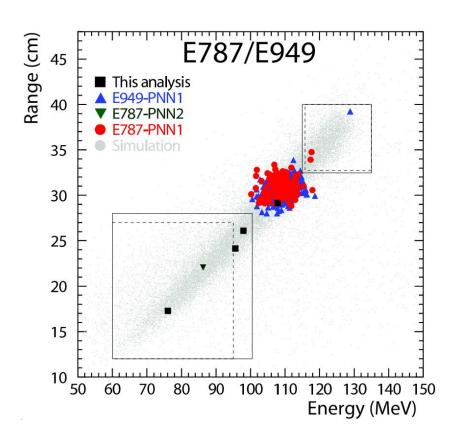


...and three new events were observed



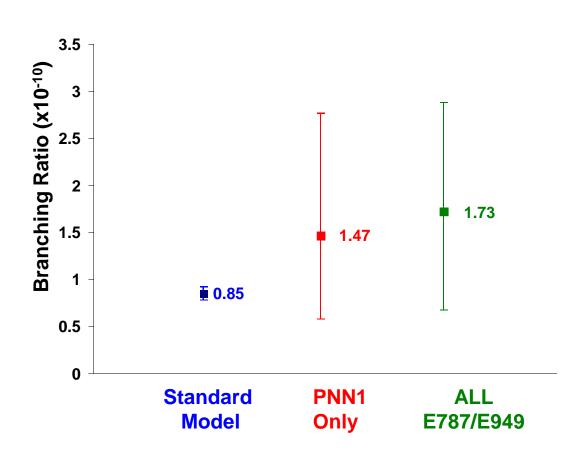
The probability that all three events were due to background only is 3.7%

The E787/E949 experiment has observed 7 candidate events in total



The probability that all seven events were due to background only is 0.1%

The combined E787/E949 branching ratio remains consistent with the Standard Model expectation



The next measurements of $K \to \pi \nu \overline{\nu}$ are just around the corner

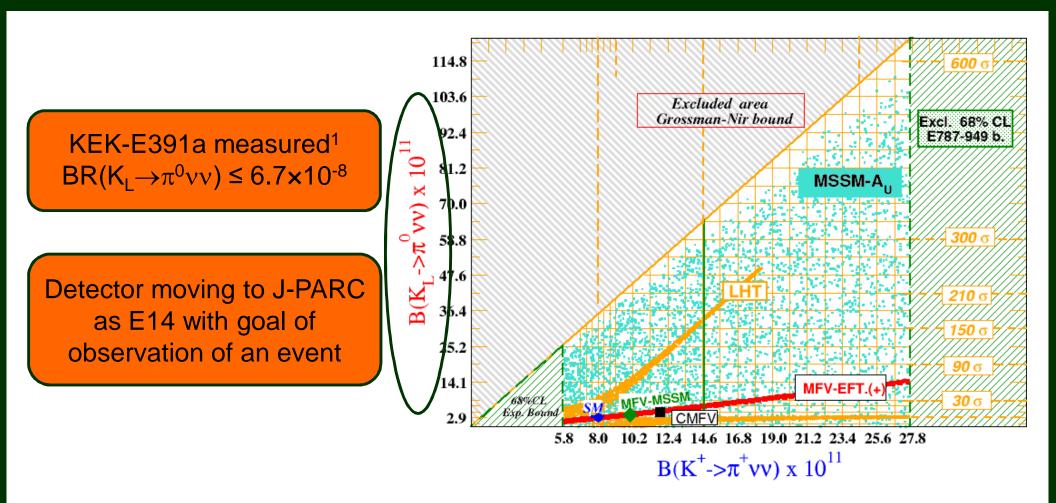


Figure: FlaviaNet - http://www.lnf.infn.it/wg/vus/

¹E391a Collaboration, J. Ahn, et al., Phys. Rev. Lett. **100** (2008) 201802.

The next measurements of $K \to \pi \nu \overline{\nu}$ are just around the corner

NA62 (formerly NA48/3) in preparation at CERN

²Goal is to measure 100 events with S/B = 10:1

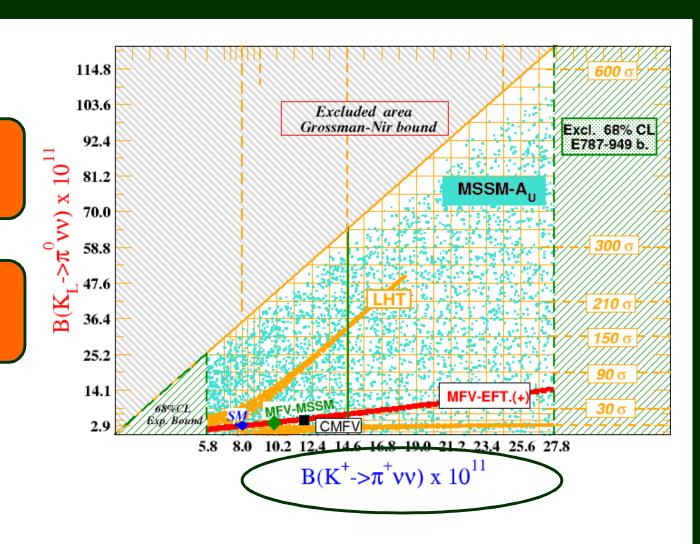
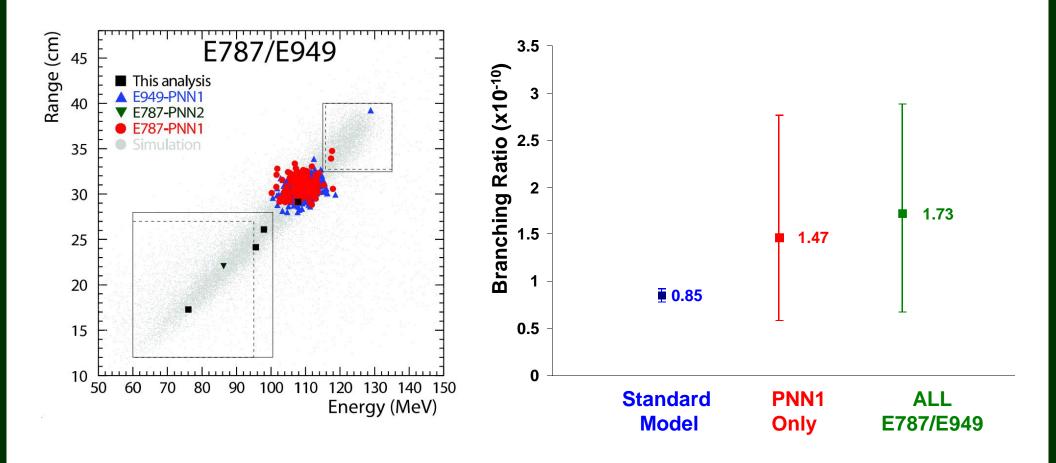


Figure: FlaviaNet - http://www.lnf.infn.it/wg/vus/

²G. Anelli et al., CERN-SPSC-2005-013, CERN-SPSC-P-326, Jun 2005.

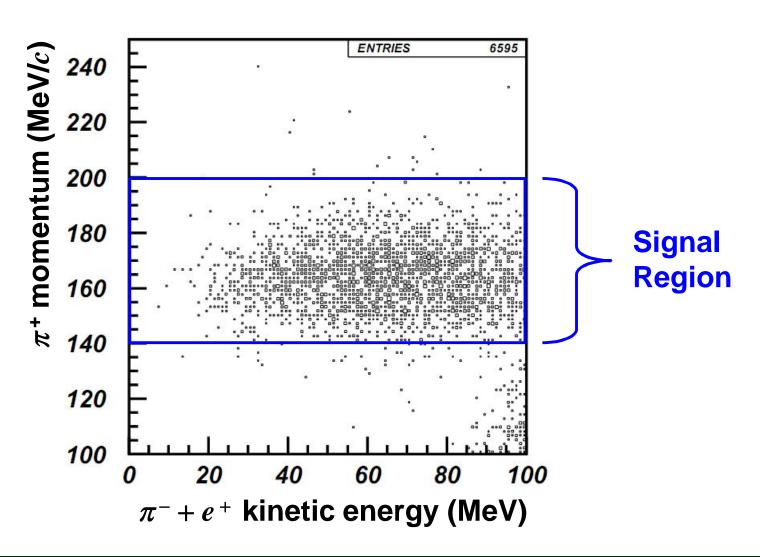
Thank you



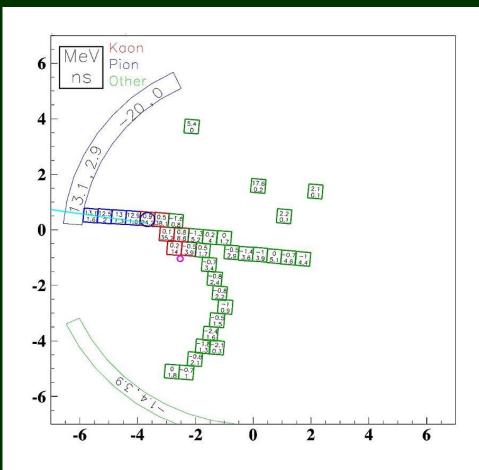
END

The $\pi^+\pi^-e^+\nu$ decay can mimic signal when the π^- and e^+ have low kinetic energies

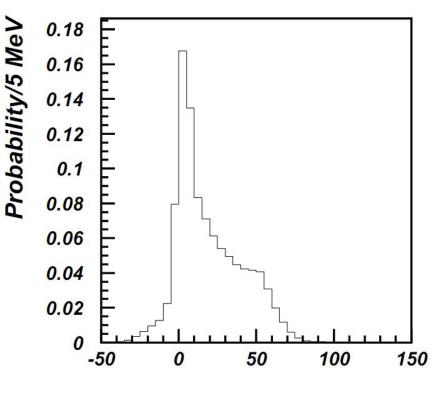
Simulated $K^+ \rightarrow \pi^+ \pi^- e^+ \nu$



Use target pattern recognition to isolate sample and estimate rejection power using simulation



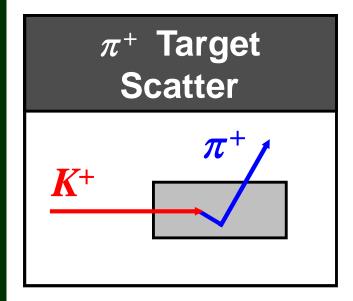
 $K^+ \to \pi^+ \, \pi^- \, e^+ v$ sample is isolated using target pattern recognition

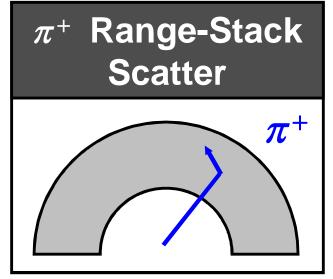


Energy of stopped pi- (MeV)

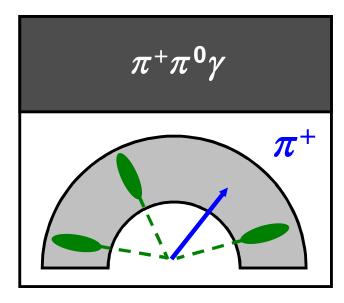
Estimate the rejection power of the target pattern recognition using simulated data supplemented by the measured π^- energy deposit in scintillator

In practice the three $\pi^+\pi^0$ backgrounds have to be disentangled

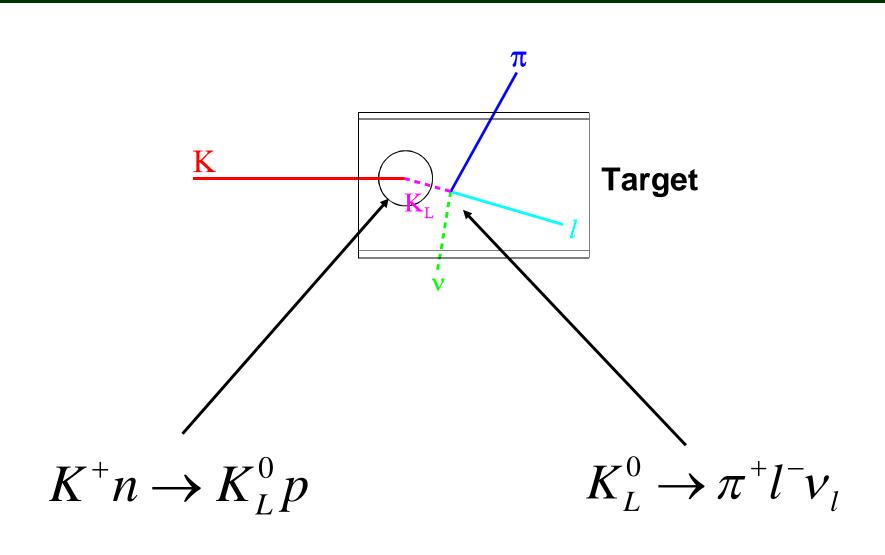




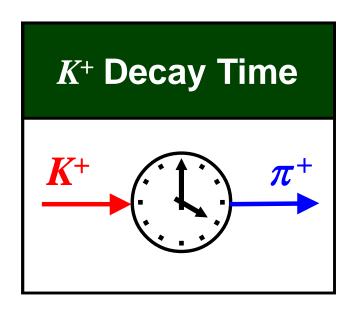
The π^+ target scatter sample made by inverting the photon veto has contamination from π^+ range-stack scatter and $K^+ \rightarrow \pi^+ \pi^0 \gamma$ processes



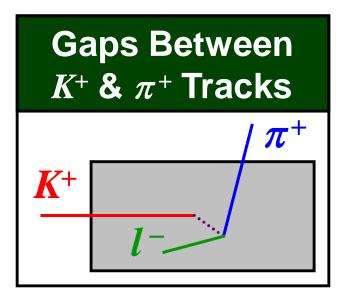
The charge exchange background is dominated by the neutral kaon decay $K_L \rightarrow \pi^+ \ l^- \ v$



$K_{\rm S}$ is suppressed by K^+ decay time condition and $K_{\rm L}$ by gaps between the K^+ and π^+ tracks



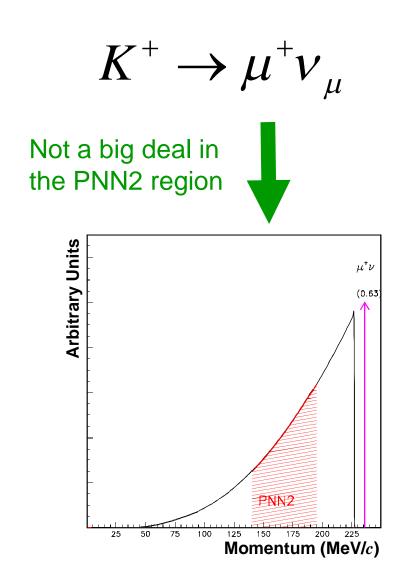
 $K_{\rm S}$ lifetime is only 0.1 ns so the K^+ decay time condition suppresses $K_{\rm S}$ decays



 $K_{\rm L}$ lifetime is 50 ns so target pattern recognition is used to identify gaps between the K^+ and π^+ tracks

Additional suppression is provided by detecting the negatively charged lepton

Due to kinematic constraints, the troublesome μ^+ decays are the multi-body ones

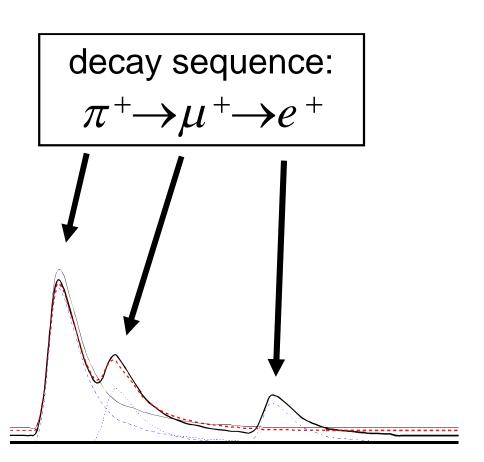


$$K^+ \to \mu^+ \nu_\mu \gamma$$

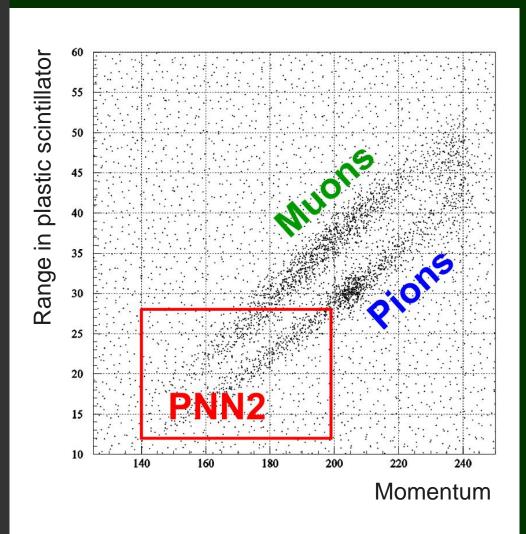


$$K^+ \rightarrow \mu^+ \pi^0 \nu_\mu$$

Muon processes are suppressed kinematically and by identification of the π^+ decay

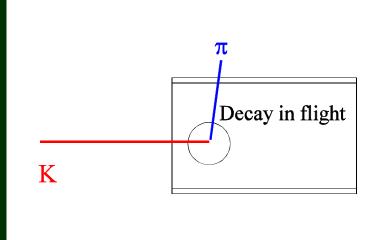


Pulse shape in the range-stack stopping counter

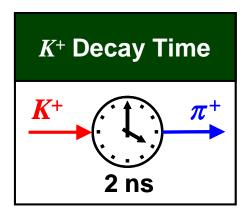


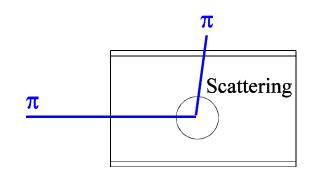
Muons have different kinematic signature in the detector than pions

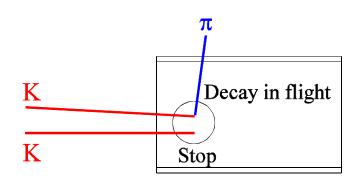
Beam backgrounds come from events other than a single kaon decaying at rest in the target



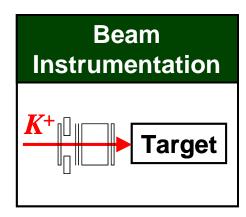


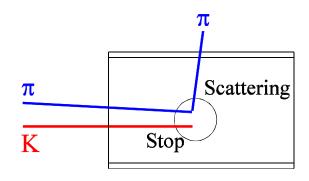




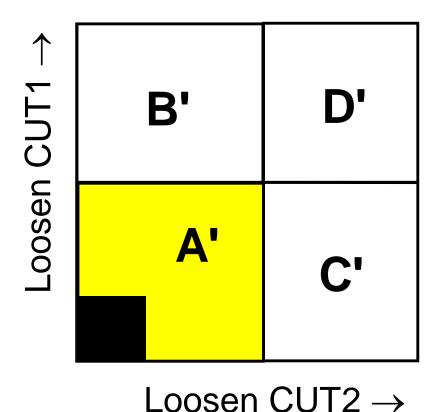


Double-beam





Outside-the-box studies compare background estimates to direct measurements just outside the signal region



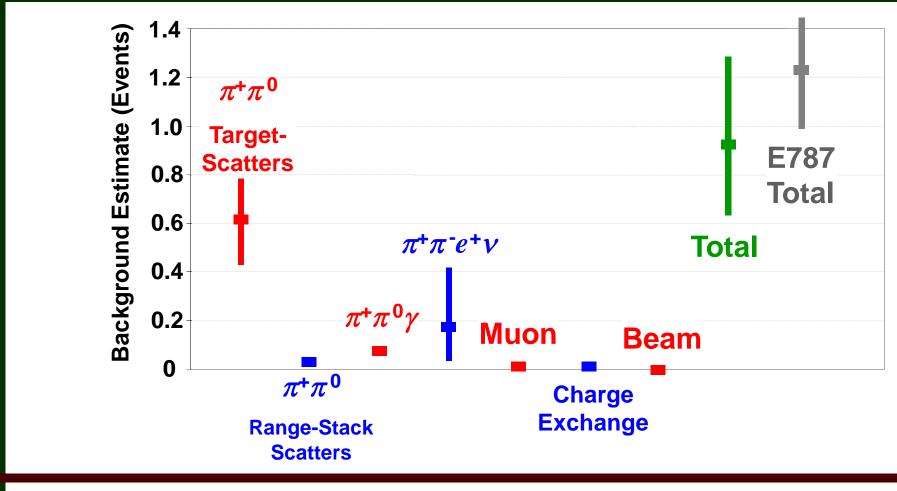
Signal region is masked out; the background estimate and direct measurement for region A' are compared directly

If CUT1 and CUT2 are uncorrelated, the two values will agree

The outside-the-box results demonstrate that the systematic uncertainties assigned to the backgrounds were reasonable

Region	$N_{\rm exp.}$	$N_{\rm obs.}$	$\mathcal{P}(N_{\mathrm{obs.}}; N_{\mathrm{exp.}})$	Combined
Target-scatter ID	$0.79^{+0.46}_{-0.51}$	0	$0.45 \ [0.29, 0.62]$	N/A
Photon Veto 1	$9.09^{+1.53}_{-1.32}$	3	$0.02 \ [0.01, 0.05]$	$0.05 \ [0.02, 0.14]$
Photon Veto 2	$32.4^{+12.3}_{-8.1}$	34	$0.61 \ [0.05, 0.98]$	0.14 [0.01, 0.40]

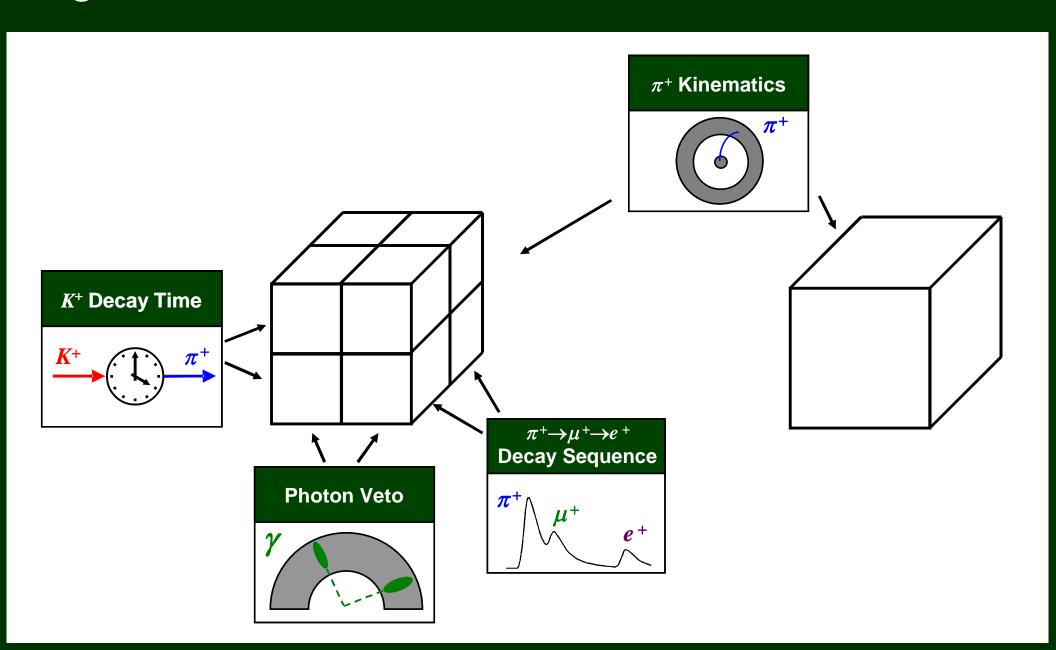
Compared to the E787-PNN2 analysis, our total background was decreased by 24% and total acceptance increased by 63%



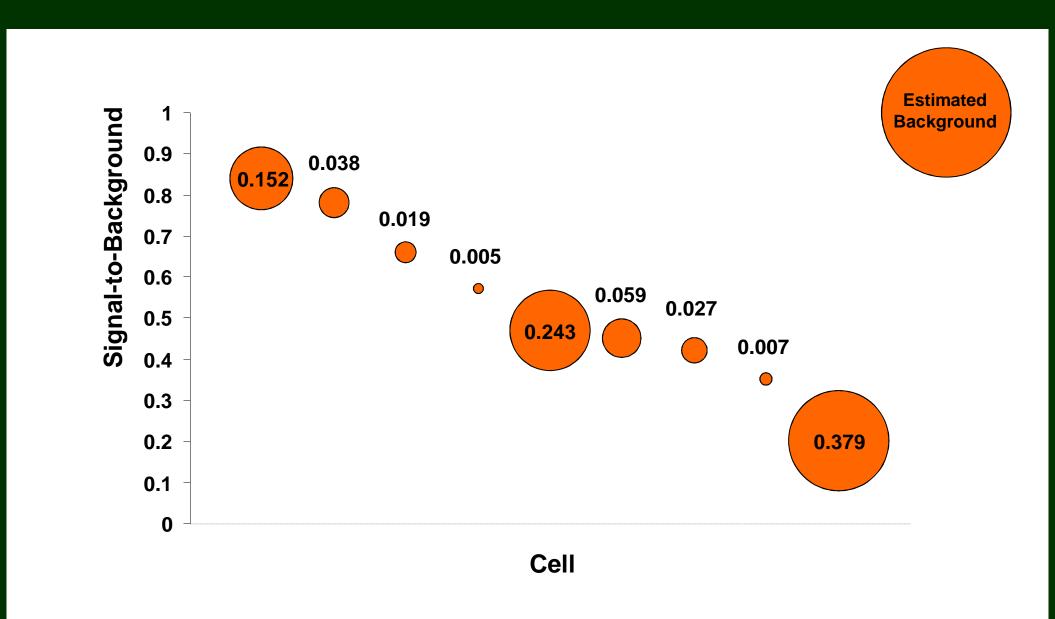
	E787-PNN2	This Analysis
Total kaons	1.73×10^{12}	1.70×10^{12}
Total acceptance	0.84×10^{-3}	1.37×10^{-3}



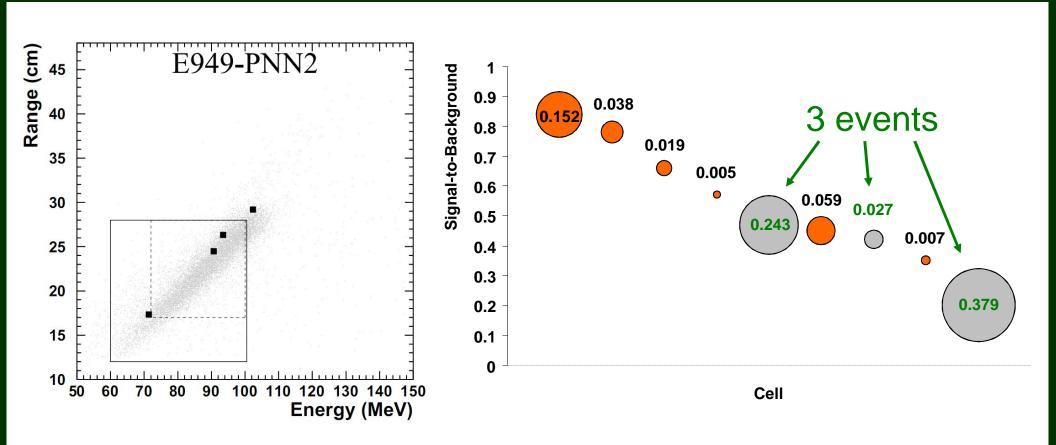
By further tightening four of the cuts, the signal region was divided into nine cells



The relative signal-to-background levels of the nine cells varied by about a factor of 4



...and three events were observed



The probability that all three events were due to background only is 3.7%

Measured branching ratio is $B(K^+ \to \pi^+ \nu \overline{\nu}) = (7.89^{+9.26}_{-5.10}) \times 10^{-10}$